

# Reliability of IBM's Public Quantum Computers

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## ABSTRACT

One of the challenges of the current ecosystem of quantum computers (QC) is the stabilization of the coherence associated with the entanglement of the states of their inner qubits. In this empirical study, we monitor the reliability of IBM's public-access QCs network on a daily basis. Each of these state-of-the-art machines has a totally different qubit association, and this entails that for a given (same) input program, they may output a different set of probabilities for the assembly of results (including both the right and the wrong ones). Although we focus on the computing structure provided by the "Big Blue" company, our survey can be easily transferred to other currently available quantum mainframes. In more detail, we probe these quantum processors with an ad hoc designed computationally demanding quaternary search algorithm. As stated, this quantum program is executed every 24 hours (for nearly 100 days) and its goal is to put to the limit the operational capacity of this novel and genuine type of equipment. Next, we perform a comparative analysis of the obtained results according to the singularities of each computer and over the total number of executions. In addition, we subsequently apply (for 50 days) an improvement filtering to perform noise mitigation on the results obtained proposed by IBM. The Yorktown 5-qubit computer reaches noise filtering of up to 33% in one day, that is, a 90% confidence level is reached in the expected results. From our continuous and long-term tests, we derive that room still exists regarding the improvement of quantum calculators in order to guarantee enough confidence in the returned outcomes.

## KEYWORDS

Error Mitigation, IBM Q, Quantum Computers, Quality Control, Quantum Search Algorithms.

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## I. INTRODUCTION

QUANTUM computing is a very promising and radically incipient area of knowledge in contrast with the current limitations of classical computing. For instance, classic transistors have a finite physical volume, and we are already approaching the 1 nm limit. Theory predicts that after surpassing this physical dimension, manipulating the flow of the electric current (without the loss of electrons) may be operationally impossible [1]. That is why, additional computational technologies may be needed to expand the ability to solve some (new) complex problems or perform extremely convoluted calculations.

According to the scientific community, quantum computing may be the solution for addressing these issues, including new frontier challenges related to machine learning and artificial intelligence [2]. However, despite all the efforts in this discipline, the challenges are still overwhelming. Perhaps, the biggest one has to do with the stability of the quantum states and the way they are organized in an interdependent way (*entangled*). This characteristic (preservation of entanglement) is known as *coherence* and obviously, critically depends on the number of qubits in a quantum computer (QC), since

the greater the number of entangled units in the quantum system, the more sensitive the system to external fluctuations. In turn, for a group of entangled qubits, there is no certainty about the exact state in which each of them is at the individual level. It is what is called *superposition of states*. Coherence in quantum computing can be defined as the conservation of the superposition state of a system over time. In a certain way, we could link this property to a loss of the individuality of each unit (qubit), to behave as a whole (hence the term coherence), and it is a requirement of the system. This property causes the system to be extremely sensitive to interferences from the environment, and as coherence can be destroyed by simple mechanical vibrations, electromagnetic disturbances, sound waves, tiny seismic temblors, or even adverse weather effects. When this takes place, the quantum wave functions collapse as if they were being measured, and the system loses its multi-state nature. This unavoidable process is known as *decoherence*. The maintenance of coherence (or avoidance of decoherence) in this type of hardware is essential for the correct implementation and execution of quantum instructions and the derivation of the expected (and accurate) results. However, in practice it has been shown that quantum decoherence can be minimized, but it is not possible, at least today, to eliminate it completely.

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Therefore, quantum error correction (QEC) is necessary in quantum computing to protect information from errors caused by decoherence and other sources of noise at the quantum level (see [3]–[6]). However, while classic error correction uses the redundancy process to counteract errors (storing the information several times in such a way that if the copies are not the same later, you can choose the option that is generally present) this possibility is no longer feasible in quantum computing according to the no-cloning theorem [7]. Even though this theorem seems to present an obstacle to formulating a theory of QEC, some alternative strategies exist. One of these has to do with the spread of qubits through highly entangled states of several neighboring physical qubits in such a way that a state inversion event could be detected without the need for consulting the exact value of the examined qubit (which would destroy the information). These qubit aggregates (making up a compound logical qubit) are resistant to errors in the final computer. Clearly, this means that if a program requires 10 qubits to run, in practice, it will need 10 logical qubits, which can be translated into hundreds or thousands of the original physical ones. These systems, called *noisy intermediate scale quantum computers* or NISQ, are expected to provide the advantages necessary to meet the required QEC [8], [9].

All errors can be corrected if the imperfections of quantum operations and measurements are below a certain threshold and the correction can be applied repeatedly [10], [11]. However, these error thresholds also depend on the details of the physical system and quantifying them requires careful analysis of both the hardware and software implementation [12].

Although some studies have addressed the quantification of these error baselines for different platforms and QC configurations (see [13]), it is a relatively new area which needs further clarification through experimentation. For this reason, in this work, the reliability of the coherence of IBM's public quantum computers has been examined. The choice of the Big Blue's network of QCs has not been arbitrary. Two factors have influenced this decision. On one hand and for several years, IBM has provided, free of charge, some of its QC infrastructure for research and study. On the other, each piece of equipment is designed differently and with a contrasting qubit number, arrangement, and entangling layout, which determined the possibility of conducting these tests in a variety of configurations for comparative purposes.

In more detail, our experiment consisted of executing for almost 100 days, 1024 times each day, the same quaternary search algorithm (described in Section IV) on 8 IBM public quantum computers. For the sake of completeness, the characteristics of IBM's public quantum processors as well as the environmental conditions in which they are designed to operate are tackled in Section III. Our results are then presented and discussed in Section VI. Finally, some conclusions are drawn in Section IX.

## II. PREVIOUS WORKS

Closely related to our work, other very recent research efforts have carried out a verification process of the reliability of IBM quantum

processors applying different study methodologies both on the number of qubits and quantum gates. The depth algorithm fragmentation method used in [14] is applied 8192 times to 20 quantum processors in a single day, showing that recomposing fragmentations significantly mitigates noise and decoherence. On the other hand, the work carried out by [15] is based on the study of non-resonant holonomic gates of 3 qubits on the resonant ones. Demonstrating that 3-qubit non-resonant holonomic gates show higher fidelity (80%) compared to resonant gates (70%). In the experiment carried out in [16], its authors focus on a single 5-qubit quantum computer to create during an evaluator scenario three models of state evolution such as inversion recovery, Ramsey and entanglement-deentanglement. They conclude that the framework of steepest entropy-ascent quantum thermodynamics (SEAQT) can be used as a basis for error mitigation schemes. In contrast, the work carried out by [17] on a 20-qubit computer shows that the application of Fourier transforms can be taken as filters that improve the oscillation patterns of the expected data.

## III. IBM'S PUBLIC QUANTUM PROCESSORS

In 2016, IBM deployed and made publicly available the first 5-qubit cloud QC. This was followed by others that were organized into families according to their number of qubits. Each family was named after a bird (Table 1). Thus, we have the 5-qubit processors, which formed the Canary family, the 16-qubit processors such as Albatross, Penguin with 20 qubits, etc. In addition, within each family, the processors are named after a city, so within the Canary family are London, Rome, Vigo, etc. Melbourne is a 16-qubit QC included in the Albatross family. These names are usually given in a personal and loving way by the specific team behind the design and assemble of each computer.

TABLE I. LIST OF NAMES AND CATEGORIES ACCORDING TO THE NUMBER OF QUBITS OF IBM COMPUTERS [18]

Category	Qubits	Processors
Canary	5	Tenerife, Yorktown, Ourense, London, Vigo, Rome, Burlington, Valencia, Santiago
Albatross	16	Melbourne
Penguin	8–16	Austin, Tokyo, Poughkeepsie, Johannesburg, Singapore, Almaden, Boeblingen
Hummingbird	+16	Raleigh

### A. Main Characteristics and Components of the IBM Q Equipment

As it is shown in Table II, the characteristics of each quantum computer in which the experiment (detailed in Section 4) has been carried-out, as well as their corresponding numbers of qubits. The type of gates that were used for the design and construction of their circuits is also included. In more detail, u1, u2, and u3 are the three parameters that allow the building of any single qubit gate and have a duration of one unit of time [19]. In addition, the error rate that each door could

TABLE II. QUBIT ERROR RATE AND THE CHARACTERISTICS OF EACH ACTIVE PUBLIC IBM Q EXPERIENCE PROCESSOR

Name	Qubits	Error Rate Door CNOT	Basic Doors	Single-qubit u2 Error Rate
Melbourne	15	2.384e-2/1.000e+0	id, u1, u2, u3, cx	4.632e-4/3.482e-2
London	5	9.411e-3/1.430e-2	u1, u2, u3, cx, id	3.369e-4/4.624e-4
Burlington	5	9.009e-3/2.075e-2	u1, u2, u3, cx, id	3.568e-4/6.144e-4
Essex	5	8.434e-3/1.406e-2	u1, u2, u3, cx, id	3.929e-4/7.155e-4
Ourense	5	6.851e-3/2.976e-2	u1, u2, u3, cx, id	2.961e-4/9.845e-4
Vigo	5	7.772e-3/1.470e-2	u1, u2, u3, cx, id	3.673e-4/8.616e-4
Yorktown	5	1.280e-2/2.203e-2	u1, u2, u3, cx, id	6.106e-4/7.950e-4

develop at the time of the measurement is also referenced. This ratio will increase as time passes if the QC is not properly calibrated.

Specifically, IBM performs these calibrations twice a day on each quantum processor and conveniently keeps the users informed so that they can take them into account when eventually launching their programs. Calibration consists of carrying out a series of experiments to obtain precise information about the physical behaviour of each qubit. The values of the parameters that characterize a qubit are different for each qubit within the processor and among different processors, and these can even vary over time. It is possible to identify the qubit's proper frequency by sweeping through a range of frequencies and observing absorption signals. The qubit's frequency is the energy difference between the ground state and the excited state. Aside from calibration, these processors must remain in specific environmental conditions. They also need a temperature close to absolute zero 0 K ( $-273.144^{\circ}\text{C}$ ) to better account for the Heisenberg's uncertainty principle [20].

IBM public quantum computer hardware uses the characteristic known as superconductivity. The materials with this property can carry electrical currents without the resistance or loss of energy under specific circumstances. From an architectural point of view, superconductors are wrapped in the form of Josephson joints [21]. These structures (Fig. 1) are formed using two sheets of aluminum, which, under normal environmental circumstances, would behave like classical electrical circuits. However, in the subatomic world, they operate as quantum gates.

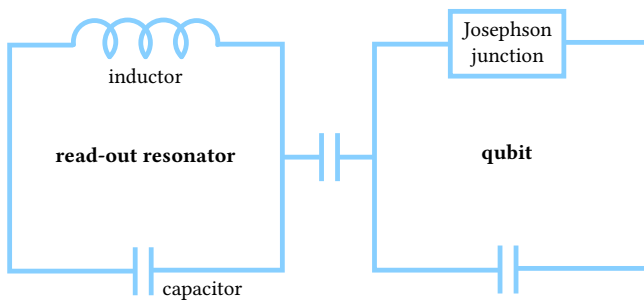


Fig. 1. Superconductivity diagram with Josephson joint (1 qubit).

The transition between the possible states of the qubits is generated by applying a certain level of energy, and because of the tunneling effect, the particle crosses the barrier (with some probability). The state of the qubit can be read by observing the energy of each aluminum sheet, which in turn, causes the decoherence of the system as a whole, but allows us to obtain information about the state in which each constituent unit (qubit) remained after the process.

### B. Quantum Computer Connectivity

As stated above, each IBM public quantum processor has a different physical architecture. Nevertheless, the logic of any given quantum

algorithm can be applied independently of the subjacent hardware. Even so, it is important to know the internal structure of these computers for several reasons. To begin with, the circuit will use, at most, all the qubits of the processor only once since the algorithms are according to the principles of their construction [22], e.g., principle 3: *long relevant decoherence times much longer than the gate operation time*). Furthermore, the application of a logic gate to several qubits requires, for greater efficiency, that they be physically interconnected in their architecture, because even if the phenomenon of entanglement itself is not conditioned by distance (two qubits could be mutually entangled even at distant points in our universe), our technological capacity to govern that entanglement to our convenience, is critically restricted by its physical separation. Therefore and a priori, the more qubits a processor has and the greater their interconnection, the better the expected results.

As can be seen in Fig. 3, Fig. 2 and Fig. 4, each IBM public QC refers to a different connectivity or type of entanglement depending on the number of qubits and their arrangement. Each circle represents a qubit, and from this possible entanglement lines emerge towards the contiguous qubits. The evolution in connectivity and lattice design is one of IBM's ongoing investigations, as shown in [18]. The debugging of errors in the gates and exposure to crosstalk is linked to the connectivity among the qubits. Therefore, the new processors are built on improvements in previous structural experiences.

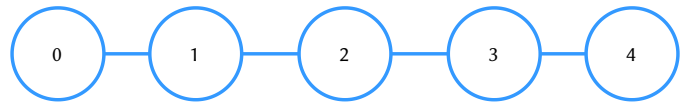


Fig. 2. Qubit arrangement for the Rome quantum computer. The structure of this processor is iterative, linking all the qubits of the processor in an orderly fashion.

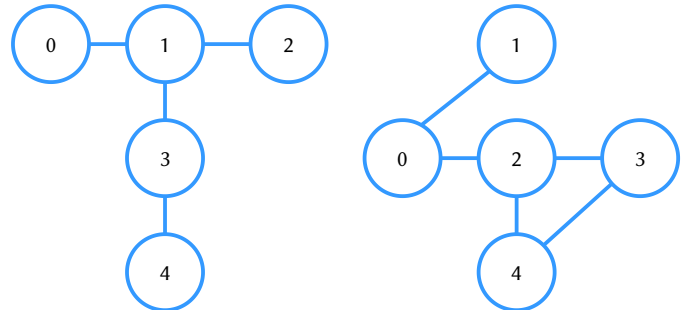


Fig. 4. Qubit arrangement for the London, Burlington, Essex, Ourense (left) and Yorktown (right) quantum computers. These architectures are a composite of Melbourne and Rome since they combine the lattice structure with the iterative one.

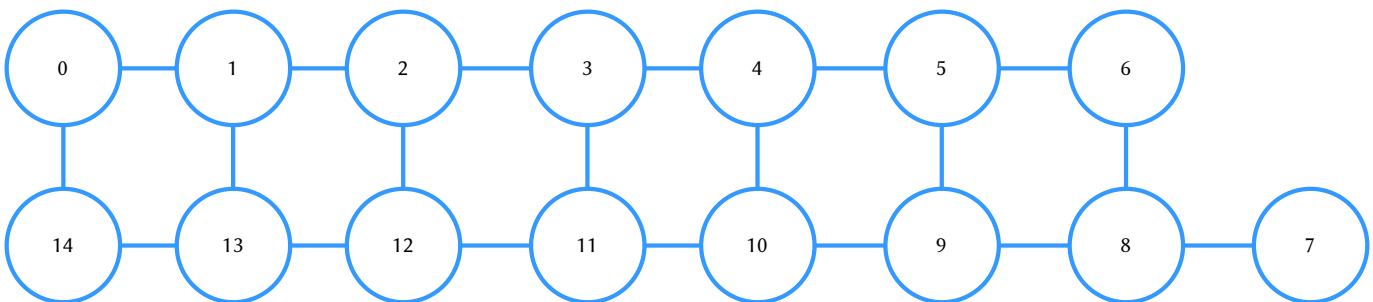


Fig. 3. Qubit arrangement for the Melbourne quantum computer. This architecture is used on processors with more than 16 qubits.

#### IV. DESIGN OF THE QUATERNARY SEARCH ALGORITHM

In classical computing, a binary tree is a data structure widely used in dynamic memory programming. Each node of the complete tree can have a left and a right child, where its complexity in the search for ordered elements in the best case is as follows:

$$(O(\log_2(N))) \quad (1)$$

where  $N$  is the number of nodes in (1). The algorithm presented here uses the data structure of a tree, but in this case, it exploits the intrinsic characteristics of the entanglement of qubits, thus managing a quaternary tree. Each node has four children. The complexity associated with searching in this structure will be considerably reduced in the best case if we transform it into a quaternary tree [23]:

$$O\left(\frac{1}{4\pi} \cdot \ln(N) - 1\right) \quad (2)$$

Besides, our search algorithm forces an iterative entanglement to verify the consistency and stability of the qubits. Therefore, if we take as a reference Grover's basic search algorithm [24], the number of iterations is equal to:

$$\Sigma U_f = (n_{\text{qubits}}) - 1 \quad (3)$$

where  $U_f$  is the so-called oracle (i.e. the unitary operator), and  $f$  is a Boolean function. According to (3), each iteration performs the addition of amplitudes until it approaches 1. As we can see,  $n_{\text{qubits}}$  will need  $n-1$  iterations to find the element of the list regardless of whether the first or last element is found. On the contrary, if we search a quaternary tree using the initial state of entanglement of two qubits  $\{00, 01, 10, 11\}$ , as shown in Fig. 5, the number of oracles to be used will be equal to the number of levels in the tree:

$$\Sigma U_f = \Sigma L \quad (4)$$

where  $L$  is the number of levels in the tree and  $U_f$  is the number of oracles in (4). Furthermore, each oracle will return the maximum possible amplitude, so unlike Grover's algorithm we only need to apply the oracle once in each iteration. In our case, we have reduced the application of oracles to two iterations. The result of each oracle is concatenated with the next one until they are finalized in a leaf of the tree. The entanglements of the qubits in each of the oracles are detailed next:

- $L1 = \{00, 01, 10, 11\}$ , where the entanglement is formulated by  $q[3]$  and  $q[2]$ . We generate a vector of states in the 4-D Hilbert space ( $H^4$ ). The expression  $q[i]$  for qubit  $q$  and  $i$  is the position of the qubit in the circuit.
- $L2 = \{0000, 0001, 0010, 0011, 0100, 0101, 0110, 0111, 1000, 1001, 1010, 1011, 1100, 1101, 1110, 1111\}$ , where the entanglement will be formulated by  $q[3]$ ,  $q[2]$ ,  $q[1]$ , and  $q[0]$ , if and only if the element has not been found in  $L1$ . We generate a vector of states in the 16-D Hilbert space.

To achieve this, first the Pauli gate X has been applied to  $q[2]$ , with the aim of changing the initial state. Then the qubits  $q[3]$  and  $q[2]$  have been entangled applying the following Hadamard gates that perform a rotation  $\pi$  around the XZ axis:

$$H|0\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \quad H|1\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle) \quad (5)$$

From here, we can write (5) as the matrix given by (6):

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \quad (6)$$

The state of entanglement of two qubits is determined by the Einstein-Podolsky-Rosen (EPR) pair [25]:

$$|q[2] q[3]\rangle = \frac{|00\rangle + |11\rangle}{\sqrt{2}} \quad (7)$$

Therefore, this state  $|q[2] q[3]\rangle$  described in (7) cannot be decomposed into pure states since no combination of complex coefficients fulfills both descriptions. Therefore, as an alternative to assembling pure states, it is possible to describe mixed states through the matrix or density operator  $\rho$ , explained in [26]. We then define the assembly of pure states as the set  $\{\rho_i | \psi\rangle$  where  $\rho_i$  are all possible states of and thus  $\Sigma \rho_i = 1$ . Then the density operator or density matrix is the result of the entanglement of several qubits:

$$\rho = \sum_i \rho_i |\psi_i\rangle\langle\psi_i| \quad (8)$$

If we write (8) in matrix form, we have:

$$\rho = \frac{1}{N} I \quad (9)$$

where  $N$  is the number of possible states in its measurement, and  $I$  is the identity matrix in (9). For example, the density matrix for a single qubit will be  $\rho = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ . In our case, with 2 qubits, the generated initial density matrices will be:

$$\{\rho | q[2] q[3]\} = \frac{1}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \quad (10)$$

If we write (10) for 4 qubits:

$$\{\rho | q[2] q[3] q[1] q[0]\} = \frac{1}{16} \begin{pmatrix} 1 & \dots & 0 \\ \vdots & 1 & \vdots \\ 0 & \dots & 1 \end{pmatrix} \quad (11)$$

where the set  $\{\rho | q[0] \dots q[n]\}$  defines the state probabilities of the entangled qubits. In our case, we want to find state 01 at level  $L1$  and then state 0111 at level  $L2$ . Hence, we change the sign of the amplitude of the states we are looking for such that:

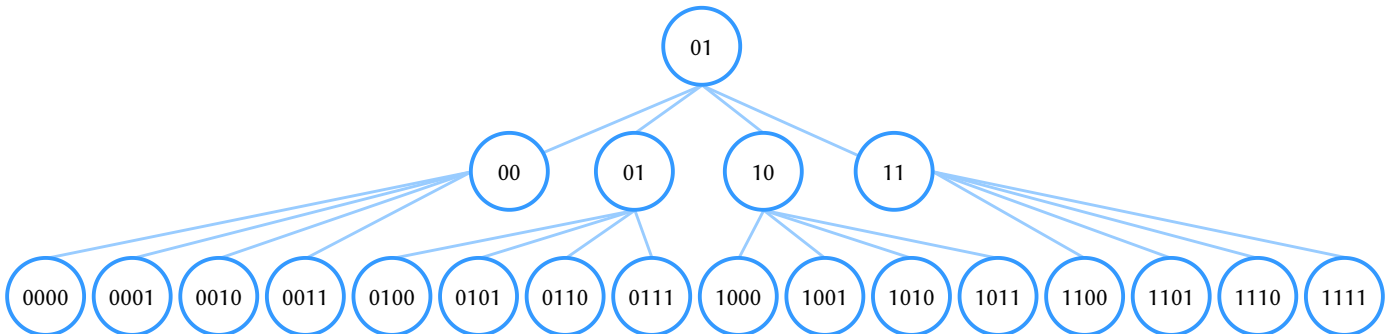


Fig. 5. Graph representation of the complete quantum quaternary tree.



$$\begin{aligned}
U_{01}|01\rangle &= -|01\rangle, \\
U_{01}|q[2]q[3]\rangle &= -|q[2]q[3]\rangle, \\
U_{0111}|0111\rangle &= -|0111\rangle, \\
U_{0111}|q[2]q[3]q[1]q[0]\rangle &= -|q[2]q[3]q[1]q[0]\rangle \\
\forall \quad q[2]q[3]q[1]q[0] \neq 0111
\end{aligned} \quad (12)$$

After applying (12), the unit transformation of the oracle  $U$  is applied. The oracles that must be applied to the algorithm are scalar according to the number of search qubits. In addition, we will only need one oracle for each level, thus reducing the Grover's algorithm:

$$L1 = (I - 2)|\omega_0\rangle\langle\omega_0| \quad L2 = (I - 2)|\omega_1\rangle\langle\omega_1| \quad (13)$$

The oracles for 2- and 4-qubit entanglement are:

$$U_{\omega_0} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \quad (14)$$

$$U_{\omega_1} = \begin{pmatrix} 0_{0000,0000} & \cdots & 0_{0000,1111} \\ \vdots & 1_{0111,0111} & \vdots \\ 0_{1111,0000} & \cdots & 0_{1111,1111} \end{pmatrix} \quad (15)$$

where  $\omega_0 = 01$  and  $\omega_1 = 0111$ . The resulting algorithm forces the qubits to iteratively intertwine during execution. In other words, the entanglement result of the first two qubits of the first level  $L1$  (given by (14) of the tree) will continue to entangle with the second level  $L2$  (given by (15)) and so on (if we expand the tree and, consequently, the number of qubits). Therefore, it should be noted that the qubits are not all initialized entangled: additions are made to the initially entangled source. One of the necessary conditions to generate a quantum circuit is the lack of breaks in the code. In this quaternary search algorithm, searching for an element of the first level  $L1$  given by (14) would be impossible. All oracles must be evaluated, and the result obtained is stored in a sheet of the second level  $L2$  given by (15).

After that, we apply Grover's star operator:

$$G_f = (2|\omega_0\rangle\langle\omega_0| - I) \quad (16)$$

Equation (16) is used to increase the amplitude of the element to be found. Recall that Grover (given by (13)) performs a search on unordered items, but in our case, they are ordered. Therefore, in principle, applying this oracle should ensure a successful outcome for the desired element. The operator defined in Grover's algorithm to increase the amplitude consists of applying the inverse of  $U_f$  in our case  $L1^{-1}$  and  $L2^{-1}$ . Considering the following assertions:

$$\begin{aligned}
L1 &= (I - 2|\omega_0\rangle\langle\omega_0|) & L2 &= (I - 2|\omega_1\rangle\langle\omega_1|) \\
L1^{-1} &= (2|\omega_0\rangle\langle\omega_0| - I) & L2^{-1} &= (2|\omega_1\rangle\langle\omega_1| - I) \\
|S_0\rangle &= |q[2]q[3]\rangle & |S_1\rangle &= |q[2]q[3]q[1]\rangle
\end{aligned} \quad (17)$$

where  $|S_0\rangle$  and  $|S_1\rangle$  in (17) are auxiliary notations for the entanglement states, then the resulting equation of the algorithm is:

$$\begin{aligned}
L1|S_0\rangle L1^{-1} + L2|S_1\rangle L2^{-1} &= \frac{1}{4\sqrt{4}} \left( (4-4) \sum_{S_0 \neq \omega_0} |S_0\rangle + 8|\omega_0\rangle \right) \\
&+ \frac{1}{16\sqrt{16}} \left( (16-4) \sum_{S_1 \neq \omega_1} |S_1\rangle + 44|\omega_1\rangle \right)
\end{aligned} \quad (18)$$

A graphical and step-by-step representation of the execution of this quaternary search algorithm (given by (18)) is shown in Fig. 6. The Qiskit framework [27] generates this graphical timeline automatically from the Python code available at GitHub. IBM has brought QC closer to the public by giving access to its processors and providing a series of intuitive tools for conducting experiments. In addition, it announced in 2020 the Quantum Educators program

which introduced training in this discipline in the classroom. To complete the teaching material, IBM offers the open-source textbook *Learn Quantum Computing Using Qiskit*. Thanks to the initiatives of the Big Blue, many students are able to train in this discipline, which would otherwise be impossible for them [28]. The code in Listing 1 shows the OpenQASM code behind the graphical representation. Both in Fig. 6 and Listing 1, we see the zero entry of the 4 qubits used defined as  $q[0]$ ,  $q[1]$ ,  $q[2]$  and  $q[3]$ . After that, the following methodology is used on the timeline:

- Steps a and b: We initialize the qubit  $q[3]$  to one, applying the Pauli gate X (U3) and then we interlace the states of the qubits  $q[2]$  and  $q[3]$  by applying the Hadamard operator (U2).
- Step c: We apply the gate CZ (Pauli Z (U1) conditioning factor) where the state we want to find is activated, in our case 01.
- Steps d, e, f, g and h: We combine the Pauli Z (U1) and X (U3) doors to perform the unitary operator and its inverse.
- Step i: We introduce the second-level qubits of the trees  $q[0]$  and  $q[1]$  with one applying the Pauli gate X (U3).
- Steps j, k, l, m and n: We combine the Pauli Z (U1) and X (U3) doors to perform the unitary operator and its inverse.
- Steps o, p, q and r: We measure the output of each qubit.

**Listing 1.** OpenQASM code of the quaternary search algorithm presented in this research work (equivalent to Fig. 6). Comments (lines beginning with #) signal the start of the carried out steps.

```

include "qelib1.inc";
qreg q[15];
creg c0[4];
barrier q[0], q[1];
barrier q[0], q[1];
barrier q[0], q[1];
barrier q[0], q[1];
barrier q[0], q[1];
barrier q[0], q[1];
barrier q[0], q[1];
barrier q[0], q[1];
#-----
a u3(3.14, 3.14, 3.14)
q[0]; u3(1.57, 3.14, 0)
q[1]; cx q[0], q[1];
#-----
b barrier q[0];
u1(3.14) q[0];
u2(0, 3.14) q[1];
u2(0, 3.14) q[2];
u3(3.14, 0, 3.14) q[3];
#-----
c cx q[2], q[3];
u3(1.57, 6.28, 3.14)
q[2]; u2(0, 3.14) q[3];
cx q[2], q[3];
u2(0, 3.14) q[2];
barrier q[2], q[3];
#-----
barrier q[2], q[3];
u1(3.14) q[2];
u2(0, 6.28) q[3];
cx q[2], q[3];
u2(0, 3.14) q[2];
barrier q[2], q[3];
#-----
p q r measure q[0] ->
c0[0]; measure q[1] ->
c0[1]; measure q[2] ->
c0[2]; measure q[3] ->
c0[3];

```

## V. DATA COLLECTION PHASE

As explained in Section I, the algorithm described above has been run on a daily basis for almost 100 days. On each execution, each of the IBM QCs (introduced in Section III.B) performed the quaternary search 1024 times. Some data is missing because of occasional maintenance/offline periods. In addition, we have considered the calibration timetable of each piece of hardware by launching our probing code both before and after this housekeeping phase. However, as we can see in Fig. 7, no improvement has been observed in the search for the desired data after the daily calibration.

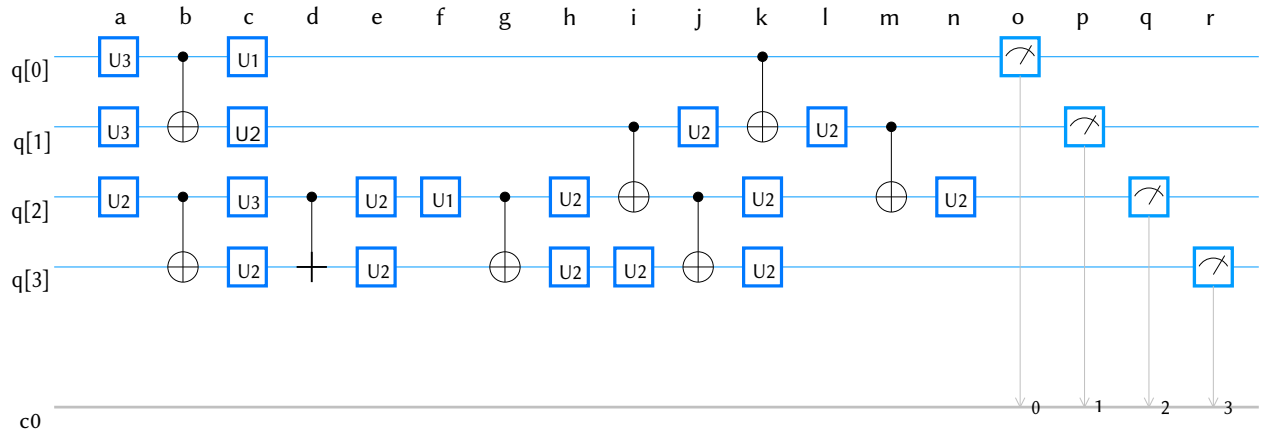


Fig. 6. Graphical representation (IBM OpenQASM 2.0 specification [27]) of the quaternary search algorithm used to probe the stability of quantum hardware. The nomenclature used by OpenQASM 2.0 indicates each turn on the qubit made by the unitary matrix or gate applied in the U form  $U(\theta, \phi, \lambda)$ ; where U2 corresponds to the Hadamard gate, U3 to Pauli X gate and U1 to Pauli Z gate.

Each run job produced, among other outcomes, a daily histogram (like the examples shown in Fig. 7 and in Fig. 8) with the probabilities of the sought result. The data of all the probabilities of the elements have been recorded to be later analyzed.

Through a graphical interface or manual code insertion with a simple Jupyter notebook written in the high-level Python language, real qubits have been used, the algorithm has been run online, and experiments have been carried out on these processors. Next, we detail the daily results of the behavior of each remote quantum computer, in addition to making a generic evaluation of all of them.

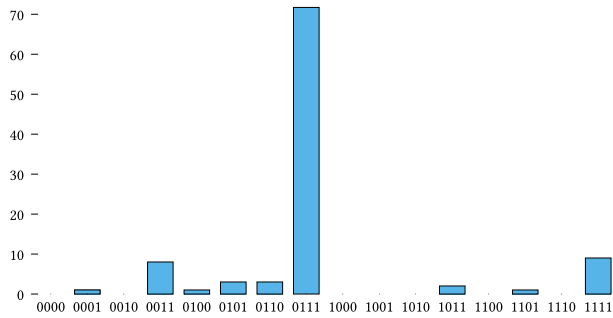


Fig. 7. Histogram of results obtained in the Yorktown QC on June 10, 2020 (after calibration). The 0111 value is the correct one.

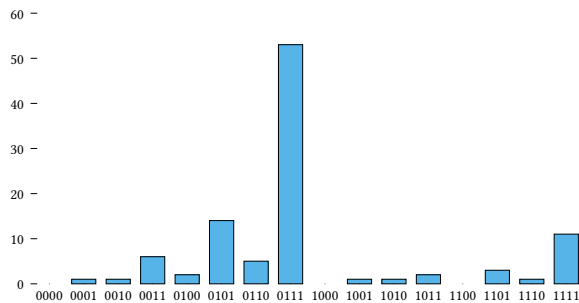


Fig. 8. Histogram of results obtained in the Burlington QC on May 29, 2020 (after calibration). The 0111 value is the correct one.

As seen in Fig. 2, Fig. 3 and Fig. 4, each computer registers a certain initial interleaving architecture. However, the execution process is supposed to generate all the necessary interleaves for any given execution. Table III shows the relationship between the original entanglements for each computer and the proposed quaternary search algorithm.

TABLE III. ENTANGLEMENT RELATIONSHIP BETWEEN THE INTRINSIC ARCHITECTURE OF EACH QUANTUM COMPUTER AND THE ENTANGLEMENT FORCED BY THE SEARCH ALGORITHM (INITIAL PROCESSOR INTERLEAVING ARCHITECTURE). HOWEVER, THERE ARE POSSIBLE PARTIAL ENTANGLEMENTS INDICATED IN THE THIRD COLUMN

	$q[3] \ \& \ q[2]$	$q[0] \ \& \ q[1]$	$(q[3] \ \& \ q[2]) \ \& \ (q[3] \ \& \ q[2])$
Melbourne	Yes	Yes	Yes
Rome	Yes	Yes	Yes
London	No	Yes	$(q[2])-(q[0]-q[1])$
Burlington	No	Yes	$(q[2])-(q[0]-q[1])$
Essex	No	Yes	$(q[2])-(q[0]-q[1])$
Ourense	No	Yes	$(q[2])-(q[0]-q[1])$
Vigo	No	Yes	$(q[2])-(q[0]-q[1])$
Yorktown	Yes	Yes	Yes

## VI. RESULTS

Next, we detail the behavior of each piece of remote hardware when running the quantum program described in Section IV. It is possible to establish two main categories: those who show a low variability and those who exhibit erratic performance over time.

### A. Quantum Computers With a Stable Trend Over Time

As shown in Fig. 9, the Yorktown computer presents the best trend to reach the probability (66%) of the correct result (0111). In addition, the difference between all the obtained values obtained does not exceed 5%. On the other hand, Fig. 10 shows the results obtained through the daily observation of the execution in this remote computer.

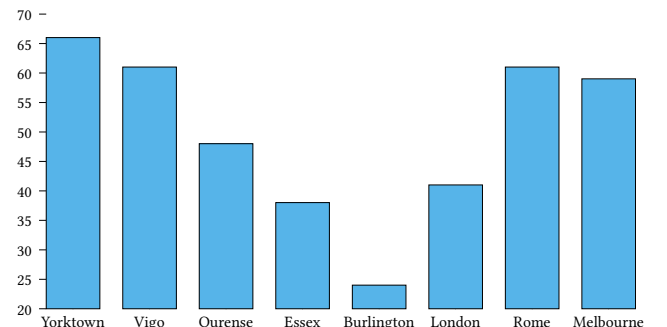


Fig. 9. Average probability of the desired result (0111) through the ~90 days of running the algorithm on several remote QC.

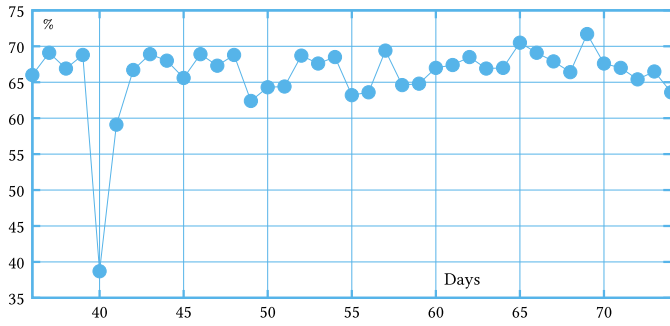


Fig. 10. Results of the Yorktown IBM equipment. Each data point refers to the probability of the desired result (0111) through the days of running the algorithm on this remote quantum computer.

### B. Unified Evaluation of All the Quantum Computers

In Fig. 9, it is shown a global (averaged over time) view of the performance of the 8 QCs. As we can see, none of them exceed the 65% average probability. Furthermore, the computer with the highest number of qubits (Melbourne) does not offer the greatest performance, which seems to indicate the importance of decoherence of entangled qubits during the process.

## VII. MITIGATION OF CROSSTALK

The term Noisy Intermediate Scale Quantum (NISQ) refers to prototype systems with 5-20 qubits that are now available for wide public use [9], as is the case of those that have been used in this work. In NISQ systems, a major source of noise such as diphyony corrupts quantum states when multiple gates or instructions are executed simultaneously [29]. Noise reduction through crosstalk mitigation in IBM quantum computers is done physically on the hardware through daily calibrations. However, this calibration is impossible for us to carry out (since we do not have privileged access to the hardware) and therefore, in this work other mitigation methods are explored through the application of software. One of the basic software tools proposed by IBM for noise mitigation in its quantum computers is the application of filters through noise matrices [30]. IBM's proposal to reduce noise is carried out through Qiskit's open `CompleteMeasFitter` class and consists of applying software filters on the initial probabilistic results. These filters are based on the creation of a noise matrix, which houses the deviations from the basic states. Therefore, any other state in superposition will be helped by a weighting in these deviations. Detailed development of this methodology can be found in the open IBM Q Experience documentation as *Measurement Error Mitigation*. The noise mitigation software has been applied to the same algorithm described in Section IV for 50 days, which was carried out, as an example, on the Santiago, Bogotá and Yorktown instruments. The reason for running this process over several days is based on the variability of daily calibrations that IBM performs at its facilities. That is, the noise matrix applied to the algorithm is different in each execution, as are the probabilistic results of the quaternary search.

The hypothesis test statistic applied for this case is Wilcoxon [31], since the test variables are adjusted to its methodology. The hypothesis proposed suggests a significant improvement in the initial results by applying noise mitigation. As can be seen in the box diagram in Fig. 11, the combined improvement exceeds 85% of the expected value. Therefore, we can conjecture that the application of noise mitigation in measurements proposed by IBM gives the expected results.

## VIII. DISCUSSION

According to the results obtained in Section VI (when compared against the expected ones), we can derive that a greater number of qubits does not guarantee a better response of a quantum computer. This is due to different reasons. In the first place, the entanglement configuration between different qubits is a determining issue. In our case, and as described in Section IV, our algorithm forces the qubits  $q[3] - q[2]$  and  $q[0] - q[1]$  to be entangled. However, only the initial interleaving architecture of the Melbourne, Rome and Yorktown computers meet this requirement, as shown in Table III and Fig. 9 (with a 60%, 62% and 66% probability, respectively). Regarding the computers London, Burlington, Essex and Ourense, where their initial interleaving configuration is the same, we verify that the results are the worst in the test, as shown in Fig. 9. This circumstance reaffirms our hypothesis about the interleaving architecture required by the tested algorithm and its direct relationship with the distilled results. In addition, we find out that not only the entanglement arrangement influences the initial structure, but also the number of qubits significantly influences it. For instance, the Melbourne computer, despite being the one with the highest number of qubits and having the required structure, it only reached a 60% probability for the expected result (also shown in Fig. 9). However, as can be seen in Fig. 7, after a hardware calibration carried out internally by IBM, the expected result improves significantly.

Finally, it has been possible to slightly balance the decoherence issue (15% improvement) by applying noise mitigation software to the results obtained (as shown in Section VII). From Fig. 11, we consider it to be a favorable result for the application of quantum computing in problem solving areas.

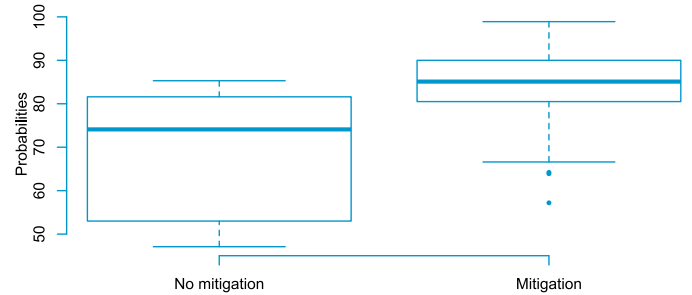


Fig. 11. Result of the application of the Wilcoxon statistic on the noise mitigation values in the quantum computers of IBM Santiago, Bogotá and Yorktown.

## IX. CONCLUSIONS

In this work, the reliability in time of a specific series of public access quantum processors has been studied through the repeated and transversal execution of the same state-of-the-art quantum algorithm. In addition, a quantum decoherence filtering proposed by IBM has been applied with a significant improvement in the results. The objective was to empirically demonstrate their current suitability for executing and a consuming resource and a computational hungry quantum program. The results obtained provide information on the probability of the correct sequences (in our case known), having shown that the results are highly dependent on the equipment in which they are executed, in turn closely related to the initial configuration of the qubits, their sequence and level of interrelation. Although for a simple task like this, it might seem if we compare with a classical scheme, that the results are not robust enough, extrapolated to tasks beyond the limits of a classical scheme, they reinforce the idea of the great potential that this technology has, even with the need to gain homogeneity over time to guarantee an adequate level of reliability in

relevant decisions, such as the business, health, or academic world. On the other hand, its suitability for research, education and the study and advancement of this technology itself has been amply demonstrated (there are currently 380,000 registered users, 1.4 trillion circuits executed and 1400 research articles published).

Future research may, to begin with, expand the number of daily executions of the algorithm. Furthermore, it could also modify the algorithm and assess the level of acceptance of qubit entanglement on the results. It would also be interesting to analyze the level of error of the quantum gates in greater detail. Finally, the fact of not knowing the calibration schedule a priori has also been a limitation.

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